

Normalisation et orthogonalité

$$\langle \psi | \psi \rangle = 1$$

$$\langle \psi_i | \psi_j \rangle = \delta_{ij}$$

Lien $\vec{\mu} / \vec{L}$:

$$\vec{\mu} = i\vec{A} \quad \text{boîte de courant}$$

$$\mu = iA \quad \text{avec } i = \frac{dq}{dt} = \frac{q}{t} = \frac{q}{\frac{2\pi\hbar}{mv}} = \frac{qv}{2\pi\hbar}$$
$$= \left(\frac{qv}{2\pi\hbar} \right) \left(\frac{\hbar^2}{mv} \right) = \frac{qv}{2} \quad \text{et } L = \frac{\hbar^2}{mv}$$

$$\text{donc } \mu = \frac{qL}{2m} \quad \text{idem vectoriellement.}$$

↙ création quantum vibration

$$\hat{H} = \hat{a}\hat{a}^\dagger + \frac{1}{2} \quad \text{↙ destruction.}$$

Oscillateur harmonique

Atome d'hydrogène dans un champ magnétique

$$\begin{aligned}
 1) \quad H &= \frac{\vec{U}^2}{2m} + V(\vec{r}) = \frac{1}{2m} [\vec{P} - e\vec{A}]^2 + V(\vec{r}) \quad \text{avec } \vec{A} = \frac{1}{2} \vec{B}_n \vec{r} \\
 &= \frac{1}{2m} (\vec{P} - \frac{e}{2} \vec{B}_n \vec{r}) (\vec{P} - \frac{e}{2} \vec{B}_n \vec{r}) + V(\vec{r}) \quad \vec{a} \cdot (\vec{b}_n \vec{c}) = \vec{b}_n \cdot (\vec{c}_n \vec{a}) \\
 &= \frac{\vec{P}^2}{2m} - \frac{e}{4m} [(\vec{B}_n \vec{r}) \cdot \vec{P} + \vec{P} \cdot (\vec{B}_n \vec{r})] + \frac{e^2}{8m} |\vec{B}_n \vec{r}|^2 + V(\vec{r}) \\
 &= \frac{\vec{P}^2}{2m} + V(\vec{r}) - \frac{e}{4m} [\vec{B}_n \cdot (\vec{r}_n \vec{P}) + \vec{B}_n \cdot (\vec{r}_n \vec{P})] + \frac{e^2}{8m} |\vec{B}_n \vec{r}|^2 \\
 &= H_0 - \frac{e}{2m} \vec{B}_n \cdot \vec{L} + \frac{e^2}{8m} |\vec{B}_n \vec{r}|^2 \\
 &\quad \downarrow \text{particule non perturbée} \quad \downarrow \text{termes quadratiques} \Rightarrow \text{négligeables devant les 2 autres} \\
 &\approx H_0 - \frac{e\hbar}{2m\hbar} \vec{B}_n \cdot \vec{L} = H_0 - \frac{\mu_B}{\hbar} \vec{L} \cdot \vec{B}
 \end{aligned}$$

$$2) \quad \text{term } -\frac{\mu_B}{\hbar} \vec{L} \cdot \vec{B} = -\frac{e}{2m} \vec{L} \cdot \vec{B} = -\vec{\pi} \cdot \vec{B} \quad \text{couplage entre m. magnétique de H (dipolaire) et le ch. magnétique} \quad \textcircled{3}$$

$$3) \quad \Psi_{nlm} \text{ sont propres de } H = H_0 - \frac{\mu_B}{\hbar} \vec{L} \cdot \vec{B} = H_0 - \frac{\mu_B B}{\hbar} L_z$$

$$H |\Psi_{nlm}\rangle = (H_0 - \frac{\mu_B B}{\hbar} L_z) |\Psi_{nlm}\rangle = H_0 |\Psi_{nlm}\rangle - \frac{\mu_B B}{\hbar} L_z |\Psi_{nlm}\rangle$$

$$\text{avec } L_z |\Psi_{nlm}\rangle = m\hbar |\Psi_{nlm}\rangle, \text{ il vient:}$$

$$H |\Psi_{nlm}\rangle = E_{nlm}^{(0)} |\Psi_{nlm}\rangle - \frac{\mu_B B}{\hbar} m\hbar |\Psi_{nlm}\rangle = (E_{nlm}^{(0)} - m\mu_B B) |\Psi_{nlm}\rangle$$

énergie propre associée
≡ motif des niveaux.

$$4) \quad E_{nlm} = E_{nlm}^{(0)} - m\mu_B \frac{B}{\hbar} = E_{nlm}^{(0)} + m \frac{\hbar}{2} \omega$$

$$5) \quad |\Psi_{100}\rangle \quad n=1 \quad l=0 \quad m=0 \quad \text{état } 1s$$

$$|\Psi_{21m}\rangle \quad n=2 \quad l=1 \quad m=0 \pm 1 \quad \text{états } 2p: |\Psi_{211}\rangle, |\Psi_{210}\rangle \text{ et } |\Psi_{21-1}\rangle$$

$$\text{Etat } |\Psi_{100}\rangle: \quad E_{100} = E_{100}^{(0)} + 0 \frac{\hbar}{2} \omega = -E_I \Rightarrow \text{origine des énergies}$$

$$\text{Etat } |\Psi_{21m}\rangle: \quad E_{21m} = E_{21m}^{(0)} + m \frac{\hbar}{2} \omega = (\frac{\hbar}{2} \Omega + E_{100}^{(0)}) + m \frac{\hbar}{2} \omega = \frac{\hbar}{2} \Omega + m \frac{\hbar}{2} \omega - E_I$$

$$|\phi_m(t=0)\rangle = \cos \alpha |\Psi_{100}\rangle + \sin \alpha |\Psi_{21m}\rangle$$

$$\begin{aligned}
 |\phi_m(t)\rangle &= \cos \alpha e^{\frac{iE_I t}{\hbar}} |\psi_{100}\rangle + \sin \alpha e^{-i(\frac{\hbar}{2}\Omega + m\frac{\hbar}{2}\omega - E_I)t} \\
 &= e^{\frac{iE_I t}{\hbar}} \left[\cos \alpha |\psi_{100}\rangle + \sin \alpha e^{-i(\Omega + m\omega)t} |\psi_{21m}\rangle \right]
 \end{aligned}$$

\uparrow $E_I = 0$ origine des énergies

$$|\phi_m(t)\rangle = \cos \alpha |\psi_{100}\rangle + \sin \alpha e^{-i(\Omega + m\omega)t} |\psi_{21m}\rangle$$

$$\begin{cases}
 x = r \sin \Theta \cos \varphi \\
 y = r \sin \Theta \sin \varphi \\
 z = r \cos \Theta
 \end{cases}$$

$$\begin{aligned}
 Y_0^0(\Theta, \varphi) &= \frac{1}{\sqrt{4\pi}} \\
 Y_1^{\pm 1} &= \pm \sqrt{\frac{3}{8\pi}} \sin \Theta e^{\pm i\varphi} \\
 &= \pm \sqrt{\frac{3}{8\pi}} \sin \Theta (\cos \varphi \pm i \sin \varphi) = \pm \sqrt{\frac{3}{8\pi}} \left(\frac{x}{r} \pm i \frac{y}{r} \right) \\
 Y_1^0 &= \sqrt{\frac{3}{4\pi}} \cos \Theta = \sqrt{\frac{3}{4\pi}} \frac{z}{r}
 \end{aligned}$$

soit en écriture d'ondensée

$$Y_1^1 = -\sqrt{\frac{3}{8\pi}} \left(\frac{x + iy}{r} \right); \quad Y_1^{-1} = \sqrt{\frac{3}{8\pi}} \left(\frac{x - iy}{r} \right)$$

$$Y_1^0 = \sqrt{\frac{3}{4\pi}} \frac{z}{r}$$

ou $x + iy = -\sqrt{\frac{8\pi}{3}} r Y_1^1$ $x - iy = \sqrt{\frac{8\pi}{3}} r Y_1^{-1}$

$$z = \sqrt{\frac{4\pi}{3}} r Y_1^0$$

ou encore finalement

$$\begin{cases}
 x = \sqrt{\frac{8\pi}{3}} r (Y_1^{-1} - Y_1^1) \\
 y = i \sqrt{\frac{8\pi}{3}} r (Y_1^1 + Y_1^{-1}) \\
 z = \sqrt{\frac{4\pi}{3}} r Y_1^0
 \end{cases}$$